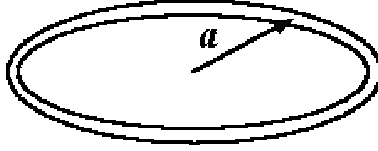


1. 3.12 The system is described by



a) From Eq. (3.106)

$$\Phi(\rho, \phi, z) = \sum_{m=0}^{\infty} \int_0^{\infty} dk e^{-kz} J_m(k\rho) [A_m(k) \sin m\phi + B_m(k) \cos m\phi]$$

where from Eq. (3.109),

$$\left. \begin{matrix} A_m(k) \\ B_m(k) \end{matrix} \right\} = \frac{k}{\pi} \int_0^{\infty} d\rho \rho \int_0^{2\pi} d\phi V(\rho, \phi) J_m(k\rho) \left\{ \begin{matrix} \sin m\phi \\ \cos m\phi \end{matrix} \right.$$

where we use $\frac{1}{2}B_0$ for $m = 0$.

b) Using cylindrical coordinates, with the origin at the center of the disc, then we have $\rho = 0$, and can use the small argument expansion for $J_m(k\rho)$

$$J_m(k\rho)|_{\rho=0} = \frac{\delta_{m0}}{\Gamma(1)} + O((k\rho)^2) = \delta_{m0}$$

$$\Phi(0, \phi, z) = \frac{1}{2} \int_0^{\infty} dk e^{-kz} B_0(k)$$

And, using Mathematica 4,

$$B_0(k) = 2kV \int_0^a d\rho \rho J_0(k\rho) = 2kV \frac{a}{k} J_1(ka) = 2VaJ_1(ka)$$

Thus, again using Mathematica 4,

$$\Phi(0, \phi, z) = Va \int_0^{\infty} dk e^{-kz} J_1(ka) = V \frac{\sqrt{z^2 + a^2} - z}{\sqrt{z^2 + a^2}} = V \left(1 - \frac{z}{\sqrt{z^2 + a^2}} \right)$$

c) We notice that for this $V(\rho, \phi)$, which is independent of ϕ , that all $A_m(k)$ vanish, and that only B_0 is nonzero. Again

$$B_0(k) = 2kV \int_0^a d\rho \rho J_0(k\rho) = 2kV \frac{a}{k} J_1(ka) = 2VaJ_1(ka)$$

$$\Phi(a, \phi, z) = Va \int_0^{\infty} dk e^{-kz} J_0(ka) J_1(ka)$$

Using Mathematic 4,

$$\int_0^{\infty} dk e^{-kz} J_0(ka) J_1(ka) = \frac{1}{2a} \left(1 - \frac{zk}{\pi a} K(k) \right)$$

where $k = \frac{2a}{\sqrt{z^2+a^2}}$, and the complete elliptic integral of the first kind is defined by

$$K(k) = \int_0^{\pi/2} \frac{d\alpha}{\sqrt{1-k^2 \sin^2 \alpha}}$$

Thus

$$\Phi(a, \phi, z) = \frac{V}{2} \left(1 - \frac{zk}{\pi a} K(k) \right)$$