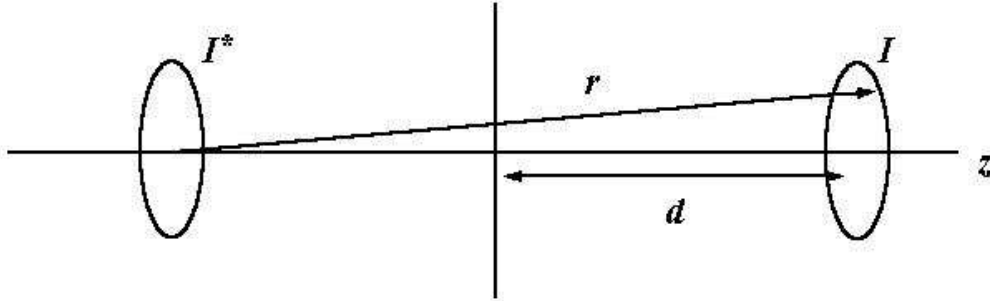


2. 5.18

a) From the results of Problem 5.17, we can replace the problem stated by the system



where  $I^*$  is equidistant from the interface and is equal to  $I^* = \frac{\mu_r - 1}{\mu_r + 1} I$ . The radius of each current loop is  $a$ . Now from Eq. (5.7)

$$\vec{F}(\text{on } I) = I \int d\vec{l} \times \vec{B}(\vec{r})$$

$$d\vec{l} \times \vec{B} = d\vec{l} \times \vec{B}_r + d\vec{l} \times \vec{B}_\theta = dl B_r (-\hat{\theta}) + dl B_\theta \hat{r}$$

By symmetry, only the  $z$  – component survives, so, from the figure

$$(d\vec{l} \times \vec{B}) \cdot \hat{z} = dl B_r \left( \frac{a}{\sqrt{4d^2 + a^2}} \right) + dl B_\theta \left( \frac{2d}{\sqrt{4d^2 + a^2}} \right)$$

So

$$F_z = \frac{2\pi a I}{\sqrt{4d^2 + a^2}} [a B_r + 2d B_\theta]$$

with  $B_r$  and  $B_\theta$  given by Eqs. (5.48) and (5.49) and  $\cos \theta = \frac{2d}{\sqrt{4d^2 + a^2}}$ ,  $r = \sqrt{4d^2 + a^2}$ , and  $I \rightarrow I^*$ .

c) To determine the limiting term, simply let  $r \rightarrow 2d$  and take the lowest non-vanishing term in the expansion of the magnetic flux density.

$$F_z = \frac{\pi a I}{d} [a B_r + 2d B_\theta]$$

$$F_z = \frac{\pi a I}{d} \left[ a \left( \frac{\mu_0 I^* a}{4d} \frac{a}{(2d)^2} \right) + 2d \left( -\frac{\mu_0 I^* a^2}{4} \left( \frac{1}{(2d)^3} \right) \right) \left( -\frac{a}{2d} \right) \right]$$

$$F_z \rightarrow -\frac{3\pi\mu_0}{32} \frac{a^4 I \times I^*}{d^4}$$

The minus sign shows the force is attractive if  $I$  and  $I^*$  are in the same direction. This same result can be gotten more directly, using

$$F_z = \nabla_z(mB_z)$$

with  $m = \pi a^2 I$ , and (from Eq. (5.64))

$$B_z = \frac{\mu_0}{4\pi} \left( \frac{2m^*}{z^3} \right)$$

with  $m^* = \pi a^2 I^*$ , and  $z = 2d$

$$F_z = \frac{\mu_0}{4\pi} 2\pi a^2 I^* \pi a^2 I \left( -\frac{3}{(2d)^4} \right) = -\frac{3\pi\mu_0}{32} \frac{a^4 I \times I^*}{d^4}$$

with agrees with out previous result.