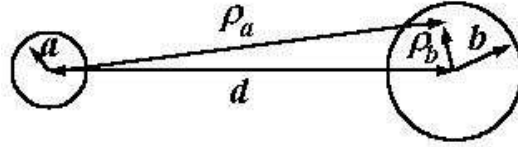


3. 5.29 The system is described by



This problem is very much like 5.26, except the wires are superconducting. We know from section 5.13 that the magnetic field within a superconductor is zero. We will be using

$$W = \frac{1}{2} \int \vec{J} \cdot \vec{A} d^3x = \frac{1}{2} \int [J_a A_{da_a} + J_b A_{da_b}]$$

Using the same arguments as applied in problem 5.26,

$$A_z = \begin{cases} -\frac{\mu I}{2\pi} \left(\ln \frac{\rho}{R} + C \right) = -\frac{\mu I}{4\pi} \left(\ln \frac{\rho^2}{R^2} + 0 \right) & \text{on the outside} \\ 0, & \text{on the inside} \end{cases}$$

Thus if we consider the second term $\frac{1}{2} \int J_b A_{da_b}$,

$$\begin{aligned} \frac{1}{2} \int J_b A_{da_b} &= \frac{1}{2} \frac{I}{\pi b^2} \int [A_{out}(\rho_a) + A_{in}(\rho_b)] \rho_b d\rho_b d\phi \\ &\simeq \frac{1}{2} \frac{I}{\pi b^2} \frac{\mu I}{4\pi} 2\pi \int_0^b \ln \frac{d^2}{a^2} \rho_b d\rho_b = \frac{1}{2} \left(\frac{\mu}{4\pi} \right) \left(2 \ln \frac{d}{a} \right) I^2 \end{aligned}$$

The first term $\frac{1}{2} \int J_a A_{da_a}$ is equal to

$$\frac{1}{2} \int J_a A_{da_a} = \frac{1}{2} \left(\frac{\mu}{4\pi} \right) \left(2 \ln \frac{d}{b} \right) I^2$$

Thus

$$W = \frac{1}{2} \left(\frac{\mu}{4\pi} \right) \left(2 \ln \frac{d^2}{ab} \right) I^2 = \frac{1}{2} \frac{L}{l} I^2$$

so

$$\frac{L}{l} = \left(\frac{\mu}{4\pi} \right) \left(2 \ln \frac{d^2}{ab} \right)$$

Now using the methods of problem 1.6, assuming the left wire has charge Q , and the right wire charge $-Q$, we find

$$\phi_{12} = \int_b^{d-a} E dr = \frac{Q}{2\pi\epsilon} \int_b^{d-a} \left(\frac{1}{r} + \frac{1}{d-r} \right) dr \simeq \frac{Q}{2\pi\epsilon} \ln \frac{d^2}{ab}$$

$$\frac{C}{l} = \frac{Q}{\phi_{12}} = \frac{2\pi\epsilon}{\ln \frac{d^2}{ab}}$$

Thus

$$\frac{L}{l} \times \frac{C}{l} = \left(\frac{\mu}{4\pi}\right) \left(2 \ln \frac{d^2}{ab}\right) \times \frac{2\pi\epsilon}{\ln \frac{d^2}{ab}} = \mu\epsilon$$