

Note that I chose the z -axis to be toward the right in the figure given in *Jackson*. Also I defined the following regions: A is the region inside the spheres; B is the region between the spheres filled with the dielectric (ϵ); C is the other region between the spheres (ϵ_0); and D is the region outside the spheres.

1 Parts a and b

I did this in full gory detail. The general expansion of Φ is

$$\Phi(\vec{x}) = \sum_{lm} Y_{lm}(\theta, \phi) [A_l r^l + B_l r^{-(l+1)}] \quad (1)$$

Or, from azimuthal symmetry,

$$\Phi(\vec{x}) = \sum_l P_l(\cos \theta) [A_l r^l + B_l r^{-(l+1)}] \quad (2)$$

This needs to be true for all $\cos \theta$. In all regions, $\Phi(\vec{x}) \neq \Phi(\cos \theta)$
 Now look at each region:

- A: $r < a$

$$\Phi(\vec{x}) = \sum_l P_l(\cos \theta) [A_l r^l] \quad (3)$$

From $\Phi(\vec{x}) \neq \Phi(\cos \theta)$, only the $l = 0$ term survives. So,

$$\Phi(\vec{x}_0) = A \quad (4)$$

- B and C: $a < r < b$

$$\Phi(\vec{x}) = \sum_l P_l(\cos \theta) [B_l r^l + C_l r^{-(l+1)}] \quad (5)$$

Again, only the $l = 0$ term survives.

$$\Phi(\vec{x}) = B + \frac{C}{r} \quad (6)$$

- D: $r > b$

$$\Phi(\vec{x}) = \sum_l P_l(\cos \theta) [D_l r^{-(l+1)}] \quad (7)$$

Again, only the $l = 0$ term survives.

$$\Phi(\vec{x}) = \frac{D}{r} \quad (8)$$

Then if we use the boundary conditions that Φ is continuous everywhere,

$$A = B + \frac{C}{a} \quad (9)$$

$$B + \frac{C}{b} = \frac{D}{b} \quad (10)$$

Now consider the \vec{D} boundary conditions:

$$\vec{E}_{(B \text{ or } C)} = -\nabla\Phi = \frac{C}{r^2} \hat{r} \quad (11)$$

$$\vec{E}_{(D)} = -\nabla\Phi = \frac{C + Bb}{r^2} \hat{r} \quad (12)$$

Then

$$\nabla \cdot \vec{D}_{(BD)} = \sigma_{free} \quad (13)$$

$$\epsilon_0 \frac{C + Bb}{b^2} - \epsilon \frac{C}{b^2} = \sigma_{BD} \quad (14)$$

$$\nabla \cdot \vec{D}_{(CD)} = \sigma_{free} \quad (15)$$

$$\epsilon_0 \frac{C + Bb}{b^2} - \epsilon_0 \frac{C}{b^2} = \sigma_{CD} \quad (16)$$

$$= \epsilon \frac{B}{b} \quad (17)$$

$$(18)$$

Or, for the inner edge/boundary,

$$\nabla \cdot \vec{D}_{(AB)} = \sigma_{AB} \quad (19)$$

$$\sigma_{AB} = \epsilon \frac{C}{a^2} - 0 \quad (20)$$

$$\nabla \cdot \vec{D}_{(AC)} = \sigma_{AC} \quad (21)$$

$$\sigma_{AC} = \epsilon_0 \frac{C}{a^2} - 0 \quad (22)$$

$$(23)$$

And the total effective charge on the inner sphere is just

$$Q = 2\pi a^2 [\sigma_{AC} + \sigma_{AB}] \quad (24)$$

$$= 2\pi a^2 \left[\epsilon \frac{C}{a^2} + \epsilon_0 \frac{C}{a^2} \right] \quad (25)$$

So the constant, C is

$$C = \frac{Q}{2\pi(\epsilon + \epsilon_0)} \quad (26)$$

And between the spheres, \vec{E} is

$$\vec{E} = \frac{Q}{2\pi r^2(\epsilon + \epsilon_0)} \hat{r} \quad (27)$$

The surface charge density induced on the inner sphere (part b) is just:

$$\sigma_{AB} = \frac{Q\epsilon}{2\pi a^2(\epsilon + \epsilon_0)} \quad (28)$$

$$\sigma_{AC} = \frac{Q\epsilon_0}{2\pi a^2(\epsilon + \epsilon_0)} \quad (29)$$

2 Part c

Now

$$\sigma_{pol} = -(\vec{P}_2 - \vec{P}_1) \cdot \hat{n}_{21} \quad (30)$$

And $P_i = (\epsilon_i - \epsilon_0)\vec{E}_i$

But $P_1 = P_A = (\epsilon_0 - \epsilon_0) \cdot 0 = 0$

Then,

$$P_0 = (\epsilon - \epsilon_0)\vec{E}_B \quad (31)$$

$$= \frac{Q}{2\pi a^2} \left(\frac{\epsilon - \epsilon_0}{\epsilon + \epsilon_0} \right) \hat{r} \quad (32)$$

And, finally, $\hat{n} = -\hat{r}$

$$\sigma_{pol} = \frac{Q}{2\pi a^2} \left(\frac{\epsilon - \epsilon_0}{\epsilon + \epsilon_0} \right) \quad (33)$$